# APPLICABILITY OF THE THREE-CONSTANT WILSON EQUATION FOR CORRELATIONS OF STRONGLY NONIDEAL SYSTEMS. I. 

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It follows from a comparison among the three-constant Redlich-Kister, Renon-Prausnitz and Wilson equations that the three-constant Wilson equation fulfills best earlier defined requirements on the normal behaviour. Further, values of correlation parameters $A_{1}, A_{2}$ corresponding to different pairs $x_{0}$ and (Gl1) $x_{0}$ at constant $C$ and related limiting activity coefficients have been calculated. Limits of sets of parameters yielding $\mathscr{G}^{\mathrm{E}}>0, \mathscr{G}^{\mathrm{E}}<0$ in the whole concentration range, an S-shaped course of $\mathscr{G}^{\mathrm{E}}\left(x_{1}\right)$ as well as an extreme on the $\log \gamma_{i}\left(x_{1}\right)$ curve have been also established. A convex course of $\left[\partial^{2} \mathscr{G} / \partial x_{1}^{2}\right]\left(x_{1}\right)$ has been proved for the case of the two-constant Wilson equation.

In previous works, the applicability of the Redlich-Kister ${ }^{1}$ (further denoted as R-K) and Renon-Prausnitz ${ }^{2}$ (NRTL) equations to correlations of systems exhibiting large positive deviations from Raoult's law ${ }^{3-5}$ has been investigated. The normal behaviour of these systems has been defined in the following manner:
l) Curve $(Q I)=\partial\left[G^{\mathrm{E}} /(2 \cdot 303 R T)\right] / \partial x_{1}$ is a monotonously decreasing function in the concentration interval ( $0 ; 1$ ), i.e.

$$
\begin{equation*}
\partial(Q I) / \partial x_{1}<0 . \tag{l}
\end{equation*}
$$

2) A less strict condition requires that the molar excess Gibbs energy be positive in the concentration interval ( $0 ; 1$ ), i.e.

$$
\begin{equation*}
\mathscr{G}^{\mathrm{E}}>0 . \tag{2}
\end{equation*}
$$

From the point of view of the course of $(G 11)\left(x_{1}\right)=\partial^{2}[G /(2 \cdot 303 R T)] / \partial x_{1}^{2}$, the following cases have been distinguished:
3) Curve (GII) $\left(x_{1}\right)$ is convex in the whole concentration range, i.e.

$$
\begin{equation*}
\partial^{2}(G I I) / \partial x_{1}^{2}=(G I I I I)>0 \tag{3}
\end{equation*}
$$

4) At a nonconvex course of $(G I I)\left(x_{1}\right)$, only one extreme (i.e. minimum) on this curve has been required.

This work has been aimed at finding for the three-parameter Wilson equation ${ }^{6,7}$ : a) how it fulfills the above given requirements on the normal behaviour, $b$ ) how the equation reflects changes in $x_{0}$ and (G11) $\left(x_{0}\right), c$ ) possible values of limiting activity
coefficients for homogeneous systems, $d$ ) its advantages and disadvantages in comparison with the earlier discussed $R-K$ and NRTL equations.

## THEORETICAL

The three-parameter equation proposed by Wilson ${ }^{6,7}$ has been used in the form

$$
\begin{equation*}
\mathscr{G}^{\mathbb{E}} /(\boldsymbol{R} T)=2 \cdot 303 Q=-C\left[x_{1} \ln \left(x_{1}+A_{1} x_{2}\right)+x_{2} \ln \left(x_{2}+A_{2} x_{1}\right)\right] \tag{4}
\end{equation*}
$$

where $A_{1}, A_{2}$ and $C$ are empirical parameters which depend only on temperature and pressure.

The ratio of activity coefficients and higher (for $n>2$ ) derivatives of $\mathscr{G}^{\mathrm{E}} / \boldsymbol{R} T$ with respect to composition are given by relations

$$
\begin{gather*}
\ln \left(\gamma_{1} / \gamma_{2}\right)=2 \cdot 303(\mathrm{Q} 1)=C\left\{\ln \left[\left(x_{2}+A_{2} x_{1}\right) /\left(x_{1}+A_{1} x_{2}\right)\right]-\right. \\
\left.-A_{2} /\left(x_{2}+A_{2} x_{1}\right)+A_{1} /\left(x_{1}+A_{1} x_{2}\right)\right\}  \tag{5}\\
\partial^{\mathrm{n}} Q / \partial x_{1}^{\mathrm{n}}=(2 \cdot 303)^{-1}(-1)^{\mathrm{n}}(n-2)!C\left\{\left[1+(n-1) A_{2} /\left(x_{2}+A_{2} x_{1}\right)\right]\right. \\
\cdot\left[\left(A_{2}-1\right) /\left(x_{2}+A_{2} x_{1}\right)\right]^{\mathrm{n}-1}-\left[1+(n-1) A_{1} /\left(x_{1}+A_{1} x_{2}\right)\right] \\
\left.\cdot\left[\left(1-A_{1}\right) /\left(x_{1}+A_{1} x_{2}\right)\right]^{\mathrm{n}-1}\right\} \tag{6}
\end{gather*}
$$

For the second derivative of the Gibbs energy with respect to composition we obtain

$$
\begin{gather*}
\partial^{2}(\mathscr{G} / R T) / \partial x_{1}^{2}=2 \cdot 303(\mathrm{G} 11)=\left(x_{1} x_{2}\right)^{-1}+ \\
+C\left\{( A _ { 2 } - 1 ) \left[\left(1+A_{2} /\left(x_{2}+A_{2} x_{1}\right)\right] /\left(x_{2}+A_{2} x_{1}\right)-\right.\right. \\
\left.-\left(1-A_{1}\right)\left[1+A_{1} /\left(x_{1}+A_{1} x_{2}\right)\right] /\left(x_{1}+A_{1} x_{2}\right)\right\} . \tag{7}
\end{gather*}
$$

Limiting activity coefficients are given by relations

$$
\begin{align*}
& \lim _{\mathrm{x}_{1} \rightarrow 0} \log \gamma_{1}^{0}=L 1=C\left(1-\ln A_{1}-A_{2}\right) / 2 \cdot 303  \tag{8}\\
& \lim _{\mathrm{x}_{1} \rightarrow 1} \log \gamma_{2}^{0}=L 2=C\left(1-\ln A_{2}-A_{1}\right) / 2 \cdot 303 \tag{9}
\end{align*}
$$

Effect of $x_{0}$ and (G11) ( $x_{0}$ )
During investigating the effect of $x_{0}$ and (G11) $\left(x_{0}\right)$, the following equations have been solved numerically

$$
\begin{gather*}
{\left[\partial^{2}(\mathscr{G} \mid \boldsymbol{R} T) / \partial x_{1}^{2}\right]_{x_{1}=x_{0}}=2 \cdot 303(\mathrm{G} 11)\left(x_{0}\right)=f_{1}\left(x_{0}, A_{1}, A_{2}, C\right)}  \tag{10}\\
{\left[\partial^{3}(\mathscr{G} \mid \boldsymbol{R} T) / \partial x_{1}^{3}\right]_{x_{1}=x_{0}}=0=f_{2}\left(x_{0}, A_{1}, A_{2}, C\right)} \tag{10a}
\end{gather*}
$$

for different values of $x_{0},(\mathrm{G} 11)\left(x_{0}\right)$ and $C$. The method of the solution has been described in an earlier work ${ }^{4}$. Results of the calculations for $C=1$ are given in Tables I-IV. The results for $C>1$ in a limited extent are at disposal in the Department of Physical Chemistry at the Institute of Chemical Technology.

The dependence of $A_{2}$ on $A_{1}$ is depicted on Fig. 1 for $(\mathrm{G} 11)\left(x_{0}\right)=0$ (critical isotherm) and for different values of $x_{0}$ and parameter $C$. Values of $A_{1}$ and $A_{2}$ lower than those indicated by the respective $C$-curve correspond to a system which splits into two phases. It would have been at $A_{1}=A_{2}=0$ for $C=1$.

The dependence of $L 1$ or $L 2$ as a function of $x_{0}$ for different $C$ is on Fig. 2 or 3, resp., for (G11) $\left(x_{0}\right)=0$ (solid curves) and (G11) $\left(x_{0}\right)=0.5^{( }$(dashed curves). It is obvious from these figures that higher values of $C$ yield lower values of limiting activity coefficients $L 1$ and $L 2$ at constant $x_{0}$ and (G11) $\left(x_{0}\right)$. If we limit ourselves


Fig. 1
Limiting Values of $A_{1}, A_{2}, C$ Securing the Condition of Thermodynamic Stability for the Wilson Equation for Different $x_{0}$ and (G11) $\left(x_{0}\right)=0$


Fig. 2
Dependence of the Logarithm of the Activity Coefficient of the First Component on $x_{0}$ and C

$$
(\mathrm{G} 11)\left(x_{0}\right)=0, \cdots(\mathrm{G} 11)\left(x_{0}\right)
$$

$=0.5$.

## Table I

Dependence of the Parameters of the Wilson Equation and Limiting Activity Coefficients on (G11) ( $x_{0}$ ) for $x_{0}=0.1,0.15$ and 0.2

| $(\mathrm{G} 11)\left(x_{0}\right)$ | $A_{1}$ | $A_{2}$ | $L 1$ | $L 2$ |
| :--- | :--- | :--- | :--- | :--- |


|  | $x_{0}=0.1$ |  |  |  |
| :--- | :---: | :---: | :---: | :---: |
|  |  |  |  |  |
| 1.00 | 0.00316 | 1.513 | 2.252 | 0.253 |
| 0.90 | 0.00476 | 1.423 | 2.135 | 0.278 |
| 0.80 | 0.00586 | 1.318 | 2.093 | 0.312 |
| 0.75 | 0.00622 | 1.266 | 2.090 | 0.329 |
| 0.70 | 0.00650 | 1.201 | 2.093 | 0.348 |
| 0.60 | 0.00684 | 1.111 | 2.115 | 0.386 |
| 0.50 | 0.00696 | 0.991 | 2.155 | 0.432 |
| 0.40 | 0.00687 | 0.894 | 2.210 | 0.482 |
| 0.30 | 0.00642 | 0.749 | 2.299 | 0.557 |
| 0.25 | 0.00611 | 0.678 | 2.352 | 0.602 |
| 0.20 | 0.00567 | 0.598 | 2.420 | 0.654 |
| 0.10 | 0.00430 | 0.412 | 2.623 | 0.827 |
| 0.05 | 0.00334 | 0.313 | 2.772 | 0.962 |

$$
x_{0}=0.15
$$

| 1.00 | 0.0108 | 1.492 | 1.751 | 0.255 |
| :--- | :--- | :--- | :--- | :--- |
| 0.90 | 0.0136 | 1.389 | 1.697 | 0.286 |
| 0.80 | 0.0155 | 1.282 | 1.686 | 0.319 |
| 0.75 | 0.0162 | 1.229 | 1.690 | 0.337 |
| 0.70 | 0.0167 | 1.175 | 1.700 | 0.357 |
| 0.60 | 0.0173 | 1.063 | 1.733 | 0.400 |
| 0.50 | 0.0174 | 0.950 | 1.782 | 0.449 |
| 0.40 | 0.0168 | 0.828 | 1.849 | 0.508 |
| 0.30 | 0.0156 | 0.699 | 1.938 | 0.583 |
| 0.25 | 0.0147 | 0.628 | 1.994 | 0.629 |
| 0.20 | 0.0135 | 0.553 | 2.065 | 0.684 |
| 0.10 | 0.0103 | 0.375 | 2.265 | 0.855 |
| 0.05 | 0.0074 | 0.259 | 2.456 | 1.020 |

$$
x_{0}=0.2
$$

| 1.00 | 0.0251 | 1.461 | 1.398 | 0.258 |
| :--- | :--- | :--- | :--- | :--- |
| 0.90 | 0.0295 | 1.348 | 1.380 | 0.292 |
| 0.80 | 0.0323 | 1.235 | 1.387 | 0.328 |
| 0.75 | 0.0336 | 1.178 | 1.399 | 0.348 |
| 0.70 | 0.0341 | 1.122 | 1.414 | 0.369 |
| 0.60 | 0.0346 | 1.007 | 1.456 | 0.416 |

Table I
(Continued)

| (G11) $\left(x_{0}\right)$ | $\boldsymbol{A}_{1}$ | $\boldsymbol{A}_{2}$ | $L 1$ | $L 2$ |
| :---: | :---: | :---: | :---: | :---: |
|  |  |  |  |  |
| 0.50 | 0.0342 | 0.892 | 1.513 | 0.469 |
| 0.40 | 0.0326 | 0.771 | 1.585 | 0.532 |
| 0.30 | 0.0298 | 0.645 | 1.679 | 0.612 |
| 0.25 | 0.0279 | 0.578 | 1.738 | 0.660 |
| 0.20 | 0.0254 | 0.505 | 1.809 | 0.719 |
| 0.10 | 0.0187 | 0.340 | 2.015 | 0.893 |
| 0.05 | 0.0134 | 0.233 | 2.205 | 1.060 |

to $C=1$, then the highest values of $L 1$ and $L 2$ in a homogeneous system may approach infinity at $(\mathrm{G} 11)\left(x_{0}\right)=0$. Even at $C>1$, the Wilson equation enables to reach higher $L 1$ values for homogeneous systems in comparison with the $\mathrm{R}-\mathrm{K}$ and NRTL equations. E.g. for $C=1 \cdot 1$ or $C=1 \cdot 2\left(x_{0}=0 \cdot 1 ;(\mathrm{G} 11)\left(x_{0}\right)=0\right)$ it follows from Fig. 2 or $3: L 1=1 \cdot 91, L 2=0.62$ or $L 1=1 \cdot 73, L 2=0 \cdot 5$, resp. On the other hand the NRTL equation yields $L 1=1 \cdot 5, L 2=0.65$ at the same conditions and at an optimum value $\alpha=0.6$. The $\mathrm{R}-\mathrm{K}$ three parameter equation fails completely at these conditions.

## Conditions for a Monotonous Course of (Q1) and an S-Shaped Course of $\mathscr{G}^{\mathrm{E}}$

The monotonous course of $(\mathrm{Q} 1)$ is closely related to the existence of an extreme on curve (Q1) ( $x_{1}$ ) or $\gamma_{\mathrm{i}}\left(x_{1}\right)$, resp. Contradictory or even erroneous conclusions in this respect are found in literature and therefore we will analyze Eq. (4) in detail. Let us investigate the limiting case, i.e.

$$
\begin{equation*}
\hat{\partial}(\mathrm{Q} 1) / \partial x_{1}=0 . \tag{11}
\end{equation*}
$$

A small rearrangement of Eq. (6) yields

$$
\begin{equation*}
\alpha x_{1}^{2}+\beta x_{1}+\gamma=0, \tag{12}
\end{equation*}
$$

where

$$
\begin{aligned}
& \alpha=1-2\left(A_{1}+A_{2}\right)+\left(A_{1}+A_{2}\right)^{2}+2 A_{1} A_{2}\left[1-\left(A_{1}+A_{2}\right)+A_{1} A_{2} / 2\right], \\
& \beta=-1+4 A_{1}-2 A_{1}^{2}+A_{1} A_{2}\left[-4+2\left(A_{1}+A_{2}\right)-A_{1} A_{2}\right], \\
& \gamma=A_{1}^{2}\left(A_{2}^{2}+1-2 / A_{1}\right) .
\end{aligned}
$$

During solving Eq. (12), the following cases may occur: A. Eq. (12) has two solutions for $x_{1} \in(0,1)$. As it is shown in Appendix 1, this case is excluded for the Wilson equation regardless of the value of parameter C . This case would correspond to three extremes on curve $\mathscr{G}^{\mathbb{E}}\left(x_{1}\right)$. B. The equation has one solution for $x_{1} \in(0,1)$. One inflexion point is on curve $\mathscr{G} \mathrm{E}\left(x_{1}\right)$ at this composition and the course need not obey condition (2). Eq. (12) may also hold in the case of a zero discriminant

$$
\begin{equation*}
D=\beta^{2}-4 \alpha \gamma=0 \tag{13}
\end{equation*}
$$

This case, however, corresponds to an inflexion point with the first derivative equal to zero and it may be inferred from Eq. (13) that this is satisfied for

$$
\begin{gather*}
A_{2}=1 / A_{1}  \tag{14a}\\
\left(A_{2}\right)_{\mathrm{t}, 11}=\left[A_{1}\left(3-2 A_{1}\right) \pm 2\left(A_{1}-1\right) \sqrt{ } A_{1}\left(A_{1}-1\right)\right]\left[A_{1}\left(4-3 A_{1}\right)\right]^{-1} \tag{14b}
\end{gather*}
$$



Fig. 3
Dependence of the Logarithm of the Activity Coefficient of the Second Component on $x_{0}$ and $C$

$$
(\mathrm{G} 11)\left(x_{0}\right)=0, \cdots(\mathrm{G} 11)\left(x_{0}\right.
$$ $=0.5$.



Fig. 4
Sets of Parameters Satisfying Conditions $1 \mathscr{G}^{\mathrm{E}}>0$ and $\partial(\mathrm{Q} 1) / \partial x_{1}<0$, possibly $\mathscr{G}^{\mathrm{E}}<0$ and $\partial(\mathrm{Q} 1) / \partial x_{1}>0$ in the whole concentration range 2 set of parameters $A_{1}, A_{2}$ displaying an extreme on the curve (Q1) $\left(x_{1}\right), 3$ set of parameters yielding an S-shaped course of $\mathscr{G}^{\mathrm{E}}\left(x_{1}\right)$.

Table II
Dependence of the Parameters of the Wilson Equation and Limiting Activity Coefficients on (G11) $\left(x_{0}\right)$ for $x_{0}=0.25,0.3$ and 0.35

| (G11) $\left(x_{0}\right)$ | $A_{1}$ | $A_{2}$ | $L 1$ | $L 2$ |
| :--- | :--- | :--- | :--- | :--- |


|  | $x_{0}=0.25$ |  |  |  |
| :--- | :---: | :---: | :---: | :---: |
|  |  |  |  |  |
| 1.00 | 0.0500 | 1.410 | 1.122 | 0.263 |
| 0.90 | 0.0558 | 1.288 | 1.128 | 0.300 |
| 0.80 | 0.0594 | 1.170 | 1.152 | 0.340 |
| 0.75 | 0.0605 | 1.111 | 1.169 | 0.362 |
| 0.70 | 0.0611 | 1.053 | 1.190 | 0.385 |
| 0.60 | 0.0611 | 0.938 | 1.240 | 0.435 |
| 0.50 | 0.0592 | 0.823 | 1.304 | 0.493 |
| 0.40 | 0.0556 | 0.707 | 1.381 | 0.561 |
| 0.30 | 0.0500 | 0.586 | 1.480 | 0.645 |
| 0.25 | 0.0464 | 0.522 | 1.541 | 0.696 |
| 0.20 | 0.0420 | 0.456 | 1.613 | 0.757 |
| 0.10 | 0.0303 | 0.305 | 1.821 | 0.937 |
| 0.05 | 0.0214 | 0.207 | 2.010 | 1.112 |

$$
x_{0}=0.3
$$

| 1.00 | 0.091 | 1.331 | 0.898 | 0.270 |
| :--- | :--- | :--- | :--- | :--- |
| 0.90 | 0.097 | 1.204 | 0.923 | 0.311 |
| 0.80 | 0.100 | 1.083 | 0.960 | 0.356 |
| 0.75 | 0.101 | 1.024 | 0.983 | 0.380 |
| 0.70 | 0.101 | 0.967 | 1.008 | 0.405 |
| 0.60 | 0.099 | 0.855 | 1.066 | 0.459 |
| 0.50 | 0.095 | 0.744 | 1.135 | 0.521 |
| 0.40 | 0.087 | 0.635 | 1.218 | 0.594 |
| 0.30 | 0.077 | 0.523 | 1.320 | 0.682 |
| 0.25 | 0.079 | 0.465 | 1.362 | 0.736 |
| 0.20 | 0.064 | 0.404 | 1.454 | 0.800 |
| 0.10 | 0.045 | 0.268 | 1.665 | 0.986 |
| 0.05 | 0.032 | 0.182 | 1.856 | 1.160 |

$$
x_{0}=0.35
$$

| 1.00 | 0.155 | 1.214 | 0.715 | 0.282 |
| :--- | :--- | :--- | :--- | :--- |
| 0.90 | 0.161 | 1.088 | 0.755 | 0.388 |
| 0.80 | 0.162 | 0.971 | 0.803 | 0.377 |
| 0.75 | 0.161 | 0.915 | 0.830 | 0.403 |
| 0.70 | 0.159 | 0.861 | 0.859 | 0.430 |
| 0.60 | 0.152 | 0.757 | 0.923 | 0.489 |

Table II
(Continued)

| 0.50 | 0.142 | 0.656 | 0.996 | 0.555 |
| :--- | :--- | :--- | :--- | :--- |
| 0.40 | 0.129 | 0.557 | 1.082 | 0.632 |
| 0.30 | 0.112 | 0.457 | 1.187 | 0.725 |
| 0.25 | 0.102 | 0.406 | 1.249 | 0.781 |
| 0.20 | 0.091 | 0.352 | 1.323 | 0.847 |
| 0.10 | 0.063 | 0.233 | 1.534 | 1.040 |
| 0.05 | 0.044 | 0.158 | 1.724 | 1.217 |

C. The equation has no solution for $x_{1} \in(0,1)$ and then curve (Q1) $\left(x_{1}\right)$ is either monotonously decreasing and $\mathscr{G}^{E}\left(x_{1}\right)>1$ or it is monotonously increasing and $\mathscr{G}^{\mathrm{E}}\left(x_{1}\right)<0$.
Let us find conditions at which the extreme may appear at concentration limits, i.e. at $x_{1}=0$ and $x_{1}=1$. If the solution should yield $x_{1}=0$, it must hold

$$
\begin{equation*}
\gamma=A_{1}^{2}\left(1+A_{2}^{2}-2 / A_{1}\right)=0 \tag{15}
\end{equation*}
$$

and for $A_{1} \neq 0$ we obtain

$$
\begin{equation*}
A_{2}=\sqrt{ }\left(2 / A_{1}-1\right) . \tag{16}
\end{equation*}
$$

This dependence is on Fig. 4 denoted as $a$.
If the solution is to yield $x_{1}=1$, it must hold on the contrary

$$
\begin{equation*}
\alpha+\beta+\gamma=-2 A_{2}+A_{2}^{2}+A_{1}^{2} A_{2}^{2}=0 \tag{17}
\end{equation*}
$$

and consequently

$$
\begin{equation*}
A_{2}=2 /\left(1+A_{1}^{2}\right) . \tag{18}
\end{equation*}
$$

This dependence is denoted as $b$ in Fig, 4 and it is also found as the dashed curve in Fig. 1. The hatched area on Fig. 4 between curves a and b represents the set of parameters $A_{1}, A_{2}$ (regardless of the value of parameter $C$ ) yielding an extreme on curve (Q1) $\left(x_{1}\right)$ or an extreme on activity coefficient curves.
As it has been stated earlier, Eq. (12) has only one solution for $x_{1} \in(0,1)$ and, consequently, curve $\mathscr{G E}\left(x_{1}\right)$ has only one inflexion point. The S-shaped course of $\mathscr{G}^{\mathrm{E}}\left(x_{1}\right)$ (with the consequent change in the sign of $\mathscr{G}^{£}$ ) may occur only for opposite signs at logarithms of limiting activity coefficients. Let us again investigate the cases
for which the limiting values of the activity coefficients equal unity. We obtain for this case from Eqs (8) and (9) (again regardless of the value of parameter C)

$$
\begin{align*}
& A_{2}=1-\ln A_{1}  \tag{19}\\
& A_{2}=\exp \left(1-A_{1}\right) .
\end{align*}
$$

## Table III

Dependence of the Parameters of the Wilson Equation and Limiting Activity Coefficients on (G11) $\left(x_{0}\right)$ for $x_{0}=0.4,0.45$ and 0.5

| $(\mathrm{Gl1})\left(x_{0}\right)$ | $A_{1}$ | $A_{2}$ | $L 1$ | $L 2$ |
| :--- | :--- | :--- | :--- | :--- |

$x_{0}=0.4$

| 1.00 | 0.256 | 1.051 | 0.370 | 0.301 |
| :--- | :--- | :--- | :--- | :--- |
| 0.90 | 0.255 | 0.937 | 0.620 | 0.352 |
| 0.80 | 0.249 | 0.833 | 0.676 | 0.405 |
| 0.75 | 0.244 | 0.784 | 0.706 | 0.433 |
| 0.70 | 0.238 | 0.737 | 0.738 | 0.463 |
| 0.60 | 0.222 | 0.647 | 0.805 | 0.526 |
| 0.50 | 0.203 | 0.561 | 0.882 | 0.597 |
| 0.40 | 0.181 | 0.465 | 0.970 | 0.678 |
| 0.30 | 0.154 | 0.391 | 1.075 | 0.775 |
| 0.25 | 0.139 | 0.347 | 1.138 | 0.833 |
| 0.20 | 0.123 | 0.301 | 1.212 | 0.902 |
| 0.10 | 0.084 | 0.199 | 1.422 | 1.098 |
| 0.05 | 0.058 | 0.135 | 1.611 | 1.279 |

$$
x_{0}=0.45
$$

| 1.00 | 0.406 | 0.841 | 0.459 | 0.332 |
| :--- | :--- | :--- | :--- | :--- |
| 0.90 | 0.389 | 0.755 | 0.516 | 0.387 |
| 0.80 | 0.367 | 0.675 | 0.576 | 0.445 |
| 0.75 | 0.354 | 0.638 | 0.608 | 0.475 |
| 0.70 | 0.341 | 0.601 | 0.641 | 0.507 |
| 0.60 | 0.311 | 0.531 | 0.710 | 0.574 |
| 0.50 | 0.279 | 0.462 | 0.788 | 0.648 |
| 0.40 | 0.244 | 0.394 | 0.876 | 0.732 |
| 0.30 | 0.205 | 0.325 | 0.981 | 0.883 |
| 0.25 | 0.184 | 0.289 | 1.044 | 0.894 |
| 0.20 | 0.161 | 0.251 | 1.117 | 0.964 |
| 0.10 | 0.109 | 0.166 | 1.325 | 1.166 |
| 0.05 | 0.074 | 0.113 | 1.513 | 1.349 |

Table IV
Dependence of the Parameters of the Wilson Equation and Limiting Activity Coefficients on (G11) $\left(x_{0}\right)$ for $x_{0}=0.5$

| (G11) $\left(x_{0}\right)$ | $A_{1}=A_{2}$ | $L I=L 2$ | (G11) ( $\left.x_{0}\right)$ | $A I=A 2$ | $L I=L 2$ |
| :---: | :---: | :---: | :---: | :---: | :---: |
|  |  |  |  |  |  |
| 1.00 | 0.611 | 0.383 | 0.40 | 0.316 | 0.797 |
| 0.90 | 0.562 | 0.440 | 0.30 | 0.262 | 0.902 |
| 0.80 | 0.513 | 0.500 | 0.25 | 0.234 | 0.963 |
| 0.75 | 0.489 | 0.532 | 0.20 | 0.204 | 1.035 |
| 0.70 | 0.465 | 0.565 | 0.10 | 0.136 | 1.241 |
| 0.60 | 0.416 | 0.634 | 0.05 | 0.093 | 1.426 |
| 0.50 | 0.351 | 0.715 |  |  |  |

These dependences are denoted as $c$ and $d$ in Fig. 4. Curve $d$ is also on Fig. 1. The hatched area between both curves represents the set of parameters corresponding to the S-shaped course of $\mathscr{G} \mathrm{E}\left(x_{1}\right)$.

A more concrete representation of the course of $\mathscr{G}^{E}$ or $(Q 1)$ resulting from the three-


Fig. 5
Sets of Limiting Activity Coefficients Yielding Courses Defined in the Caption to Fig. 4


Fig. 6
Values of $L 1$ and $L 2$ corresponding to pairs $A_{1}, A_{2}$ for $C=1$
-parameter Wilson equation may be obtained from Fig. 5. This figure contains dependences from Fig. 4 transformed into variables $L 1$ and $L 2$.

Fig. 5 reveals some limitations of the Wilson equation pointed at by Orye and Prausnitz ${ }^{8}$. The Wilson equation cannot correlate correctly systems exhibiting an extreme on curve $(\mathrm{Q} 1)\left(x_{1}\right)$ or $\gamma_{i}\left(x_{1}\right)$ as long as both limiting activity coefficients are greater than $1.359[=\exp (1-\ln 2)$.$] . Once, however, one activity coefficient is$ lower than this value (and simultaneously greater than unity), the second one may be arbitrarily large. For activity coefficients lower than unity may the extreme on curve $\gamma_{i}\left(x_{1}\right)$ occur theoretically at arbitrarily low activity coefficients.

On turning back to strongly nonideal systems, the extreme on curve (Q1) ( $x_{1}$ ) may according to the Wilson equation appear, as it is obvious from Fig. 1, only for highly unsymmetrical systems and it is more probable at higher values of (G11) ( $x_{0}$ ) and $C$. The dependence of $L 1$ and $L 2$ on $A_{1}$ and $A_{2}$ is on Fig. 6 for $C=1$.

## Conditions Imposed on the Course of $(\mathrm{G} 11)\left(x_{1}\right)$

The normal behaviour of (G11) ( $x_{1}$ ) has been defined by inequality (3). From Eq. (6) we obtain

$$
\begin{align*}
(\mathrm{G} 1111) & =2 / x_{1}^{3}-2 C\left(1-A_{1}\right)^{3}\left(x_{1}+A_{1} x_{2}\right)^{-3}-6 C A_{1}\left(1-A_{1}\right)^{3}\left(x_{1}+A_{1} x_{2}\right)^{-4}+ \\
& +2 / x_{2}^{3}-2 C\left(1-A_{2}\right)^{3}\left(x_{2}+A_{2} x_{1}\right)^{-3}-6 C A_{2}\left(1-A_{2}\right)^{3}\left(x_{2}+A_{2} x_{1}\right)^{-4} \tag{20}
\end{align*}
$$

It is shown in Appendix II that condition (3) is always satisfied for $C=1$. The course of (G11) $\left(x_{1}\right)$ as given by Eq. (4) is always convex for this value of $C$. The values of (G1111) in unsymmetrical systems at $C>1$ become negative in a definite concentration range, namely for $x_{0}=0.2$ at $C>12$; for $x_{0}=0.15$ at $C>1.33$; for $x_{0}=0.1$ at $C>1.05$. These values should serve only for orientation and they depend to a certain extent (very little) on values of (G11) $\left(x_{0}\right)$. The occurrence of more extremes on curves (G11) $\left(x_{1}\right)$ is little probable and none were observed in the studied range of $x_{0}$, (G11) $\left(x_{0}\right)$ and $C$.

## CONCLUSIONS

We have calculated parameters $A_{1}, A_{2}, C$ in the Wilson three-constant equation, which correspond to systems with different values of $x_{0}$ and (G11) $\left(x_{0}\right)$. Besides that, the corresponding values of the limiting activity coefficients have been calculated. In contrast to the $\mathrm{R}-\mathrm{K}$ and NRTL equations, the limiting activity coefficients may reach even infinitely large values. This is one of the reasons which brings about
good results if this equation is used for homogeneous systems. This favourable property is supported by the fact that the equation with values of $C \leqq 1$ cannot reproduce the behaviour of systems which are heterogeneous in a certain concentration range. High values of the limiting activity coefficients provided by the $\mathrm{R}-\mathrm{K}$ and NRTL equations lead to splitting into two liquid phases.

Considerable attention has been paid to the course of $\mathscr{G}^{\mathrm{E}}\left(x_{1}\right)$ and $\partial \mathscr{G}^{\mathrm{E}} / \partial x_{1}$. The analysis has led to the conclusion that curve $\mathscr{G}^{E}\left(x_{1}\right)$ can have only two extremes, i.e. it can describe an S -shaped course of $\mathscr{G}^{\mathrm{E}}\left(x_{1}\right)$. Only one inflexion point may appear on curve $\mathscr{G}^{\mathrm{E}}\left(x_{1}\right)$ and this corresponds to an extreme on curves $\gamma_{\mathrm{i}}\left(x_{1}\right)$. A set of parameters has been found (Fig. 4) which yields an extreme on curves $\gamma_{i}\left(x_{i}\right)$ or an S-shaped behaviour of $\mathscr{G}^{\mathrm{E}}\left(x_{1}\right)$. This set has been transformed and sets of limiting activity coefficients have been found which correspond to $\mathscr{G}^{E}>0 ; \mathscr{G}^{\mathrm{E}}<0$, an S-shaped behaviour of $\mathscr{G} E\left(x_{1}\right)$ or possibly an extreme on curve $\gamma_{i}\left(x_{1}\right)$. At the same time it has been found out that at simultaneous occurrence of an extreme on curves $\gamma_{i}\left(x_{1}\right)$ are the limiting activity coefficients bounded above by a value of 1.355 without being bounded below. It has been proved for the Wilson two-constant equation that it always yields a convex course of (G11) ( $x_{1}$ ).

## APPENDIX I

Let us investigate whether the quadratic equation (12) can have two roots on the interval ( 0,1 ). It would have to hold in this case

$$
\begin{equation*}
0<\frac{-\beta \pm \sqrt{ }\left(\beta^{2}-4 \alpha \gamma\right)}{2 \alpha}<1 . \tag{A-I}
\end{equation*}
$$

As the value of $\alpha$ is always positive, it may be proved by an analysis of relation $(A-I)$ that it holds if and only if the following conditions are satisfied simultaneously:
1)
2)
3)
4)
5)

$$
\left.\begin{array}{ll}
\beta^{2}-4 \alpha \gamma & \geq 0  \tag{A-2}\\
2 \alpha+\beta & >0 \\
\alpha+\beta+\gamma & >0 \\
\beta & <0 \\
\gamma & >0
\end{array}\right\}
$$

It may be shown that for $A_{1}<1$ or $A_{1}>1$ are conditions 4,5 or 2,3 , respectively, excluded.

## APPENDIX II

Let us consider at the beginning the first three terms of Eq. (20)

$$
\begin{equation*}
F_{1}\left(x_{1}\right)=2 / x_{1}^{3}-2 C\left(1-A_{1}\right)^{3}\left(x_{1}+A_{1} x_{2}\right)^{-3}-6 C A_{1}\left(1-A_{1}\right)^{3}\left(x_{1}+A_{1} x_{2}\right)^{-4} \tag{A-3}
\end{equation*}
$$

Let us try to find such composition (at non-negative parameters $A_{1}, A_{2}, C$ ) at which this expression would be equal to zero. After rearranging we obtain

$$
\begin{gather*}
A_{1}^{4}+4 A_{1}^{3}\left(1-A_{1}\right) x_{1}+6 A_{1}^{2} x_{1}^{2}\left(1-A_{1}\right)^{2}+4 A_{1}\left(1-A_{1}\right)^{3} x_{1}^{3}(1-C)+ \\
+x_{1}^{4}\left(1-A_{1}\right)^{4}(1-C)=0 \tag{A-4}
\end{gather*}
$$

This relation simplifies considerably for $C=1$. Eq. ( $A-4$ ) reduces in this case to (for $A_{1}>0$ )

$$
\begin{equation*}
A_{1}^{2}+4 A_{1}\left(1-A_{1}\right) x_{1}+6 A_{1}^{2} x_{1}^{2}\left(1-A_{1}\right)^{2}=0 \tag{A-5}
\end{equation*}
$$

The discriminant of this equation is equal to

$$
\begin{equation*}
D=-8 A_{1}^{2}\left(1-A_{1}\right)^{2}, \tag{A-6}
\end{equation*}
$$

it is always negative and consequently no real $x_{1}$ exists for which would Eq. ( $A-4$ ) hold and, followingly, $F_{1}\left(x_{1}\right)$ can never be equal to zero for $C=1$. The same results are obtained after similar considerations performed with second three terms of Eq. (19).

Two further negative terms remain in Eq. ( $A-4$ ) at $C>1$ and $F_{1}\left(x_{1}\right)$ may be then not only equal to zero but it can also assume negative values, which was in fact observed.

## LIST OF SYMBOLS

$A_{1}, A_{2}, C$ constants in the Wilson equation
$\mathscr{G}_{\mathrm{E}}^{\mathrm{E}} \quad$ molar excess Gibbs free energy
(G11) ((G111), (G1111)) second (third, fourth) derivative of the molar Gibbs free energy with respect to composition divided by $2 \cdot 303 R T$
LI, L2 limiting values of the decadic logarithm of the activity coeficient of the $i$-th component
(G11) ( $x_{0}$ ) ordinate of the minimum on the curve (G11) ( $x_{1}$ )
$Q=\mathscr{G}^{\mathrm{E}} /(2 \cdot 303 R T)$ dimensionless molar excess Gibbs free energy
(Q1), (Q11) first, second derivative of $Q$ with respect to composition
$R \quad$ universal gas constant
$T \quad$ absolute temperature
$x_{i} \quad$ mole fraction of the $i$-th component
$\gamma_{i} \quad$ activity coefficient of the $i$-th component
$\gamma_{\mathrm{i}}^{0} \quad$ limiting activity coefficient of the $i$-th component
$\alpha, \beta, \gamma \quad$ parameters in Eq. (12)

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